مدرسازی اثر ناخالصی نوع P بر نویز شدت نسبی درلیسرای نقطه کوانتمی

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چکیده- در این مقاله به مدرسازی اثر ناخالصی نوع P بر مشخصه نویز شدت نسبی در لیزرهای نقطه کوانتمی در طول موج 1.3 میکرومتر پرداخته شده است. معادلات نرخ الکترونها و حفره‌ها در سه تراز گسته‌شده در نقاط کوانتمی به همراه گستره در تراز بهره‌وری در چاه کوانتمی و ناحیه سد پتاسیل به همراه معادله نرخ فوتون‌ها در رزیم سیگنال کوچک حول نقطه کار به فرم خطی در اینجا شده و سپس منابع نویز لجیسیمی اضافه گردیده است. محاسبات نشان می‌دهد که ناخالصی نوع P باعث کاهش میزان نویز شدت نسبی در لیزرهای نقطه کوانتمی می‌گردد. با افزایش جریان جریان، نور خروجی لیزر نقطه کوانتمی آلاییده شده افزایش یافته و از میزان نویز شدت نسبی کاسته می‌شود. این رفتار کاهشی نویز در اثر افزایش جریان و مقادیر نماد شدت نویز محاسبه شده با تابع تجربی گزارش شده همگامی دارد.

کلیدواژه- لیزر نقطه کوانتمی، ناخالصی نوع P، گین غیرخطی، منابع نویز، نویز شدت نسبی.

Theoretical Modeling of p-doping Effect on Relative Intensity Noise in 1.3µm InAs/GaAs Quantum Dot Lasers

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Abstract- In this paper we investigate theoretically the effect of p-doping on the relative intensity noise characteristics of 1.3µm InAs/GaAs quantum dot lasers. The rate equations for electrons and holes in three QDs' discrete levels, wetting and barrier states besides the rate of photons have been linearized in small signal regime around the operating point. Calculations demonstrate that p-doping would decrease the relative intensity noise of the quantum dot lasers. By increasing the injection current, the output power rises and the noise intensity reduces in p-doped QD laser which is in agreement with reported experimental results.

Keywords: Quantum dot lasers, P-doping, Nonlinear gain, Langevin sources, Relative intensity noise.
Theoretical Modeling of p-doping Effect on Relative Intensity Noise in 1.3µm InAs/GaAs Quantum Dot Lasers

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1 Introduction

In recent decade, semiconductor quantum dot (QD) lasers have been widely attracted to realize enhanced performances demanded for high speed optical communication systems [1]. From theoretical point of view, three dimensional confinement of carriers in quantum dots would results in the formation of the population inversion under much lower current injection. Since the successful report on growth of self-assembled InAs/GaAs(InP) QDs in Stranski-Krastanow (SK) mode with considerable crystal quality and high areal density, intensive experimental efforts have been devoted to realize proposed superior performances of QD lasers in the range of 1.3(1.55)µm wavelength [1]. Although low threshold current at room temperature have been observed in some real 1.3 µm InAs/GaAs QD lasers, temperature stability of these lasers had degraded from theoretical predictions [2]. This is attributed to the closely energy spaced of the discrete hole states in QDs. One effective suggested method to overcome thermal broadening of holes is to provide excess hole concentrations via moderate p-doping in GaAs barrier. Higher modal gain and more temperature stability of threshold current have been observed in p-doped QD lasers experimentally [3]. However injection of more excess holes to the wetting layer, leads to higher non-radiative Auger recombination [4]. One of the essential characteristics of a laser to be investigated is its noise aspect. Actually the value of the relative intensity noise (RIN) plays a key role in estimating the transmission data rate and the maximum possible distance of fibre-optic links. How much lower be the RIN value of a laser, higher transmission rate through longer distance could be reliable [5]. The first measurements on intensity noise of p-doped QD lasers in 1.3µm wavelength have been reported in 2012 by G. Lin et. al.[6]. Their results have shown a level of -135 to -155 (dB/Hz) over a wide range of frequency up to 10(GHz), for a p-doped 1.3µm quantum dot laser in which the QD layers have been grown with uniform barrier thickness [6]. In this paper we seek to study the effects of p-doping on RIN value of InAs/GaAs QD lasers from a theoretical point of view. The rest of this paper is organized as follows. In Section 2, we present the governing equations and the theoretical modeling approach. In Section 3, the results on the intensity noise analysis of a QD laser are discussed. Finally, conclusions are presented in Section 4.

2 Theoretical Modeling

Regarding the various system of rate equations for modelling characteristics of QD lasers, we develop a rigorous theory based on the theoretical model of Ref. [7]. For evaluating the relative intensity noise of 1.3µm QD lasers, the separate electron and hole states for GaAs-barrier and wetting layer have been included in addition to the nonlinearity of gain. The definition and value of all the parameters used in this model have been outlined in table I. Since the concentration of carriers in wetting layer state, strongly affects the Auger recombination process, the carrier relaxation rate have been calculated considering phonon-assisted and Auger relaxation [7].

Table I- Value of structural parameters

<table>
<thead>
<tr>
<th>Symbol</th>
<th>Meaning</th>
<th>Value</th>
</tr>
</thead>
<tbody>
<tr>
<td>n</td>
<td>Number of QD layers</td>
<td>8</td>
</tr>
<tr>
<td>ρ</td>
<td>Surface density of QDs(cm^-2)</td>
<td>6×10^10</td>
</tr>
<tr>
<td>b</td>
<td>Barrier thickness(nm)</td>
<td>324</td>
</tr>
<tr>
<td>----</td>
<td>----------------------</td>
<td>-----</td>
</tr>
<tr>
<td>N_e</td>
<td>Doping of barrier(cm⁻³)</td>
<td>5×10⁻⁶</td>
</tr>
<tr>
<td>L_e</td>
<td>Cavity Length(μm)</td>
<td>300</td>
</tr>
<tr>
<td>S_a</td>
<td>Area of active region(μm²)</td>
<td>2×300</td>
</tr>
<tr>
<td>θ_d</td>
<td>Coverage of QDs</td>
<td>0.157</td>
</tr>
<tr>
<td>d</td>
<td>Distance from doped layer to cap layer(nm)</td>
<td>10</td>
</tr>
<tr>
<td>τ_{bd}</td>
<td>Transition time from doped layer to cap layer(ps)</td>
<td>5.22</td>
</tr>
</tbody>
</table>

The rate equations have been written considering the lasing action from the QDs’ ground state and first excited state as the first and the second lasing modes respectively under pumping current density of \( J \) (A/m²). To investigate the RIN, we consider a time dependent small signal perturbations, \( \delta N_B^{(h)}(t) \), \( \delta f^{(h)}(t) \) and \( \delta s_j(t) \) for all the carrier concentration in barrier, densities of occupation probabilities of QD levels and photon densities around the steady state work point as follows.

\[
\dot{N}_B^{(h)}(t) = N_B^{(h)} + \delta N_B^{(h)}(t) \tag{1}
\]

\[
\dot{f}^{(h)}(t) = F_c^{(h)} + \delta f^{(h)}(t) \tag{2}
\]

\[
s_j(t) = S_j + \delta s_j(t) \tag{3}
\]

While \( N_B^{(h)}(t) \) and \( f^{(h)}(t) \) stand for the whole time dependent value of the electrons (holes) concentration in GaAs barrier, density of occupation probabilities for electrons (holes) in wetting layer and three discrete energy levels in QDs respectively and \( S_0(t) \) represents the total number of photons of the first (second) lasing mode. The steady state values of the mentioned variables, \( N_B^{(h)}(t) \), \( f^{(h)}(t) \), and \( s_j \) under a fixed injection current, would be calculated by solving twelve coupled rate equations. As well the small signal parts of the rate equations might be linearized by neglecting the quadratic and higher powers of \( \delta s \). Langevin noise sources are added to the rate equations. In rate equation of the holes in wetting layer state, the term \( \frac{N_B^{(h)} \eta_d}{\tau_{bd}} \) is added to represent the excess hole concentration in the p-doped QD lasers. It has been allowed the gain to vary with photon numbers for nonlinear gain effects [7]. Under Markovian assumptions, the Langevin noise sources might satisfy the following general relations [8].

\[
\langle F_i(t) \rangle = 0 \tag{4}
\]

\[
\langle F_i(t), F_k(t') \rangle = 2D_{kk'} \delta(t-t') \tag{5}
\]

Where angle brackets\( \langle \cdot \rangle \), show the ensemble average and \( D_{kk'} \) denotes the diffusion coefficients associated with the corresponding noise sources. The auto-correlation \( D_{kk'} \) can be considered as the terms in the rate equations which have been added to or subtracted from the number of carriers and photons [8]. The cross-correlation coefficients \( D_{kk'} \) can be interpreted as the terms which have been subtracted from the number of carriers and photons [8]. By taking Fourier transform from the linear small signal rate equations and representing them in a single matrix equation, we obtain a closed from. The intensity noise at a given frequency \( \omega \) would be characterized by the relative intensity noise (RIN) per unit bandwidth defined as [8]:

\[
\frac{\Delta f}{S_f^2} = \frac{S_0(\omega)}{S_f} \tag{6}
\]

while \( S_f = S_0 + S_1 \) stands for the total number of photons at steady state work point. The total spectrum of photons might be evaluated by solving the matrix equation.

In the next section, the numerical calculated results of RIN in the 1.3μm InAs QD lasers based on the above model would be presented.

3 Results and Discussion

In order to investigate the RIN properties of 1.3μm QD lasers, twelve rate equations have been solved under constant injection current \( J \) within Runge-Kutta method. The Light-Current (L-I) characteristics of un-doped and p-doped QD lasers have been calculated at temperature \( T=300K \) and have shown in Figure 1. As depicted in p-doped QD laser, the output power is higher than un-doped one because of excess holes which provide higher modal gain. The values of RIN in p-doped and un-doped QD lasers have been evaluated in Figure 2, for injection current density of \( J=12 \) (kA/cm²). It is demonstrated that noise intensity has lower level in p-doped QD lasers comparing with un-doped one. It can be deduced that p-doping introduces two new terms added to cross and auto-correlation coefficients of hole density at wetting layer state. Among these two terms, the term added to cross correlation coefficient has negative value with larger amplitude comparing with positive value of the new term in auto correlation coefficient. Hence RIN in p-doped QD laser gets smaller value than in un-doped one. It is also depicted that although p-doping improves noise intensity characteristics but it
reduces modulation bandwidth by decreasing resonance frequency.

![Figure 1](image1.png)

**Figure 1** - L-I characteristics of un-doped and p-doped QD lasers at temperature T=300K.

![Figure 2](image2.png)

**Figure 2** - RIN in p-doped with doping concentration of \(N_e=5\times10^{16} \text{(cm}^{-3}\text{)}\) and un-doped QD lasers at I=12 (mA)

To investigate the effect of injection current on noise intensity of p-doped QD lasers, we have evaluated RIN value of p-doped QD laser with 8 QD layers and QD areal density of \(\rho=6\times10^{10} \text{(cm}^{-2}\text{)}\) similar to structural parameters considered at [6]. Figure 3, shows RIN value of p-doped QD laser with doping concentration of \(N_a=5\times10^{16} \text{(cm}^{-3}\text{)}\) at three current density of \(J_1=6\), \(J_2=12\) and \(J_3=18\) (kA/cm²). As it is shown by increase of current and output power, the RIN value reduces from -130 to -140 (dB/Hz). It is evident that increase of output power results in decrease of relative intensity noise. This reduction behaviour due to increase of output power and calculated amplitudes of RINs are in agreement with experimental results reported recently.

![Figure 3](image3.png)

**Figure 3** - RIN value of p-doped QD laser with doping concentration of \(N_a=5\times10^{16} \text{(cm}^{-3}\text{)}\) at three current density of \(J_1=6\), \(J_2=12\) and \(J_3=18\) (kA/cm²)

### References


