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چکیده - در این مقاله با در نظر گرفتن ضریب ضریب ضایعاتی و دوره تناوب، فشردگی سالیتونهای نوری مرتثه دو در فیبرهای نوری غیرخطی بررسی شده است. در شیب سازی‌های عدیدی که از سالیتونهای نوری تاریک به دلیل پایداری بیشتر آنها در مقابل نویز نسبت به سالیتونهای نوری روشن استفاده شده است.

کلمات کلیدی - سالیتونهای روان، سالیتونهای تاریک، فیبر نوری غیرخطی، فشردگی پالس.

The Pulse Compression of Second-Order Dark Solitons

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Abstract- In this paper with considering dispersion coefficient and period normal length, pulse compression 2-order solitons in nonlinear fiber optics is investigated. In our numerical simulation we have used optical dark solitons because of its more stability in presence of noise as compared optical bright solitons.

Keywords: Bright Solitons, Dark Solitons, Nonlinear Fiber Optics, Pulse Compression.
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1 Introduction

An important application of nonlinear fiber optics consists of compressing optical pulses. Pulses shorter than 5 fs have been produced by using dispersive and nonlinear effects occurring in optical fibers [1]. Pulse compressors based on nonlinear fiber optics can be classified into two broad categories: gratings and soliton-effect compressors. A soliton-effect pulse compressor consists of only a piece of fiber whose wavelength is suitably chosen. Pulse compression takes place as a result of interplay between group velocity dispersion (GVD) and self-phase modulation (SPM). This method has been used since 1983 [2,3].

Because dark solitons are more stable in the presence of noise and spread more slowly in the presence of loss as compared with bright solitons in optical communication systems [4-6]. In the recent work, by using soliton effects, make a new method to compress dark solitons in fiber. This method has significant advantages compared with other pulse compression methods because of its low cost and simplicity. We consider dark solitons with \(N=2\) and investigate dark soliton pulse compression.

2 Second-Order Optical Dark soliton

Nonlinear Schrödinger Equation (NLSE) manages pulse propagation as follow [4]:

\[ i \frac{\partial u}{\partial \xi} - \frac{\beta_2}{2} \frac{\partial^2 u}{\partial \tau^2} + \gamma |u|^2 u = 0 \]  \hspace{1cm} (1)

where, \(u\) is the normal slowly varying amplitude of the pulse envelope, \(\xi\) is the normal distance, \(\tau\) is the normal time, \(\beta_2\) is the second-order dispersion coefficient which is a positive number for dark soliton, and \(\gamma\) is the Kerr coefficient. The soliton order is defined by:

\[ N = \sqrt{\frac{L_D}{L_{NL}}} \]  \hspace{1cm} (2)

where \(L_D\) is the dispersion length and \(L_{NL}\) is the nonlinear length of the soliton propagation. In this paper, specifically, we consider the fiber is a lossless media and the order of soliton is \(N=2\).

As known, dark soliton is the valley pulse in the uniform background. We want to simulate dark soliton pulse propagation in two cascaded fibers with Crank-Nicholson method. At the first, we lunch the pulse with the form of \(\tanh(x)\) as an input pulse into the first fiber. Figure 1 shows output pulse of dark soliton with \(N=2\) in the fiber. As shown clearly in Figure 2, during through the length of fibers, the output pulse is narrower and most intense and also, the output pulse has two gray soliton is produced.

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2.1 Second-Order Dark Soliton Period

In the Figure 3, the pulse evolution in different length is shown. As shown in Figure 3, dark soliton like the bright soliton, evolution to its first shape and repeated after characteristic length with little different in shape. There is considerable notice that if the dispersion coefficient is increased, the periodic length is decreased and the soliton pulse is constructed in lower distances. In Figures 4, two dark solitons with different dispersion coefficient are compared. If the dispersion coefficient increases, the total normal length decrease and the pulse in lower distance repeated.

2.2 The Compression of Second-Order Dark Optical Soliton In Nonlinear optical Fiber

In 2-order dark soliton, if we decrease the value of second order dispersion coefficient, the pulse compression is increased but this pulse compression has a low limit where lower that the compression of pulse is not change. In this section considering two various values for $\beta$ and
simulated numerically dark soliton propagation. In Figure 5, we take $\beta_2 = -0.4$ and we have shown dark solitons propagation in $z/z_0 = 6$ and $12$. As shown, in both $z/z_0 = 6$ and $12$, pulse is compressed in time zero. In Figure 5(a) the rate of value of input soliton pulse compression to value of output soliton pulse compression is approximately 3.5 and in Figure 5(b) this ratio is nearly 4. Also in $z/z_0 = 6$ pulse is compressed but its intensity is the same with input pulse intensity.

In Figure 6, we take $\beta_2 = -0.7$ and we have shown dark solitons propagation in $z/z_0 = 6$ and $12$. In Figure 6(a) the rate of value of input soliton pulse compression to value of output soliton pulse compression is approximately 2 and in Figure 6(b) this ratio is nearly 2.5.

![Figure 5: Pulse shapes and its compression with $\beta_2 = -0.4$.](image)

![Figure 6: Pulse shapes and its compression with $\beta_2 = -0.7$.](image)

### Conclusion

In this paper pulse compression for second-order of dark solitons in nonlinear fiber optics is investigated. In 2-order dark soliton, if we decrease the value of second order dispersion coefficient, the pulse compression is increased. Also the larger the dispersion coefficient of 2-order dark soliton is the more compression of central soliton.

### References


